

## ARTICLES

**Formation of  $\text{CH}_3\text{CFCl}^+$  from Photoionization of  $\text{CH}_3\text{CFCl}_2$ : An Application of Threshold Photoelectron Photoion Coincidence (TPEPICO) Technique****Su-Yu Chiang\****Synchrotron Radiation Research Center, No. 1, R&D Road VI, Hsinchu Science-Based Industrial Park, Hsinchu 300, Taiwan***Yu-Chang Lee and Yuan-Pern Lee***Department of Chemistry, National Tsing Hua University, 101, Sec. 2, Kuang Fu Road, Hsinchu 30013, Taiwan**Received: August 8, 2000*

The dissociation of energy-selected  $\text{CH}_3\text{CFCl}_2^+$  to form  $\text{CH}_3\text{CFCl}^+$  was studied with a molecular-beam/threshold-photoelectron-photoion-coincidence system using synchrotron radiation for ionization. Maximal releases of kinetic energy in the channel  $\text{CH}_3\text{CFCl}^+ + \text{Cl}$  at five photon energies were determined from half-widths of  $\text{CH}_3\text{CFCl}^+$  time-of-flight peaks in coincidence mass spectra. A thermochemical threshold for this channel was determined to be  $11.10 \pm 0.09$  eV after taking kinetic energy releases into account. With this threshold and heats of formation of Cl and  $\text{CH}_3\text{CFCl}_2$ , we derived  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+) = 149.0 \pm 2.1$  kcal mol<sup>-1</sup>. Ab initio calculations were also performed to predict heats of formation of  $\text{CH}_3\text{CFCl}_2$ ,  $\text{CH}_3\text{CFCl}$ , and their corresponding cations. Our experimental results agree satisfactorily with calculations. On the basis of experimental and theoretical results, we construct an energy diagram for dissociative photoionization of  $\text{CH}_3\text{CFCl}_2$  to form  $\text{CH}_3\text{CFCl}^+ + \text{Cl}$ .

**1. Introduction**

$\text{CH}_3\text{CFCl}_2$  (HCFC-141b) is used widely to replace  $\text{CFCl}_3$  (CFC-11) and  $\text{CFCl}_2\text{CF}_2\text{Cl}$  (CFC-113) in industrial applications.<sup>1,2</sup> With increased concern over its impact on atmospheric chemistry, investigations on  $\text{CH}_3\text{CFCl}_2$  become important. There are numerous reports on ultraviolet (UV) absorption cross sections<sup>3–5</sup> and photodissociation of  $\text{CH}_3\text{CFCl}_2$ ,<sup>6,7</sup> as well as reaction kinetics of  $\text{CH}_3\text{CFCl}_2$  with OH,<sup>3,8–10</sup> O(<sup>1</sup>D),<sup>11</sup> and Cl,<sup>12–16</sup> but little is known about thermochemical properties of  $\text{CH}_3\text{CFCl}_2$  and its related photoproducts, despite their importance in relation to the chemical reactivity of  $\text{CH}_3\text{CFCl}_2$ .

We measured the threshold photoelectron spectrum (TPES) and photoionization mass spectra (PIMS) of  $\text{CH}_3\text{CFCl}_2$ .<sup>17</sup> Valence ionic states were assigned on comparison of experimental vertical ionization energies in the TPES with theoretically predicted values. Three major fragments,  $\text{CH}_3\text{CFCl}^+$ ,  $\text{CH}_2\text{CCl}^+$ , and  $\text{CH}_2\text{CF}^+$ , were observed in the PIMS experiment; their respective appearance energies (AE) were determined from the onsets of the photoionization efficiency curves. The parent ion  $\text{CH}_3\text{CFCl}_2^+$  is absent. The AE value of  $11.33 \pm 0.02$  eV for production of  $\text{CH}_3\text{CFCl}^+$  is nearly identical to the nominal ionization threshold of  $11.28 \pm 0.02$  eV for  $\text{CH}_3\text{CFCl}_2$  in the TPES. Lack of parent ion and an onset for production of  $\text{CH}_3\text{CFCl}^+$  similar to that for ionization of  $\text{CH}_3\text{CFCl}_2$  imply that an unstable  $\text{CH}_3\text{CFCl}_2^+$  dissociates to form  $\text{CH}_3\text{CFCl}^+$ . More-

over, substantial kinetic energy is expected to be released on the dissociative photoionization process.

Kinetic energy released after dissociative photoionization is not taken into account in determination of AE from PIMS experiments; the resulting AE thus represents an upper limit.<sup>18,19</sup> Heats of formation reported previously<sup>17</sup> for three major fragments observed after photoionization of  $\text{CH}_3\text{CFCl}_2$  suffer from such uncertainty. In analysis of threshold photoelectron photoion coincidence (TPEPICO) spectra, we can take release of kinetic energy into account and thereby determine accurate AE of dissociative photoionization channels.<sup>20–22</sup> Recently, Evans et al.<sup>22</sup> directly determined a thermochemical threshold on dissociative photoionization in the process  $\text{SF}_6 + h\nu \rightarrow \text{SF}_5^+ + \text{F} + \text{e}^-$  with high-resolution TPEPICO, and they reported an accurate value of  $\Delta_f H_0^\circ(\text{SF}_5^+)$  that is consistent with theoretical predictions.<sup>23,24</sup> Lack of  $\text{SF}_6^+$  upon photoionization of  $\text{SF}_6$  is a behavior similar to that of  $\text{CH}_3\text{CFCl}_2$ .

In our investigation of dissociative photoionization of  $\text{CH}_3\text{CFCl}_2$ , TPEPICO mass spectra of  $\text{CH}_3\text{CFCl}_2$  were recorded at various photon energies near the appearance onset of  $\text{CH}_3\text{CFCl}^+$ . We determined a thermochemical threshold for the dissociative photoionization channel  $\text{CH}_3\text{CFCl}^+ + \text{Cl}$  to derive  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+)$  and a maximal release of kinetic energy at the appearance onset of  $\text{CH}_3\text{CFCl}^+$ . We also performed ab initio calculations to predict enthalpy changes of three isodesmic reactions, ionization energies of  $\text{CH}_3\text{CFCl}_2$  and  $\text{CH}_3\text{CFCl}$ , and accordingly predicted heats of formation of  $\text{CH}_3\text{CFCl}_2$ ,  $\text{CH}_3\text{CFCl}$ , and their corresponding cations. On the basis of experi-

\* Corresponding author. Tel: (886)3-578-0281, ext. 7315. Fax: (886)3-578-9816. E-mail: schiang@src.gov.tw.

mental and theoretical results, we discuss the  $\text{CH}_3\text{CFCl}^+ + \text{Cl}$  channel in photoionization of  $\text{CH}_3\text{CFCl}_2$ .

## 2. Experiment and Calculations

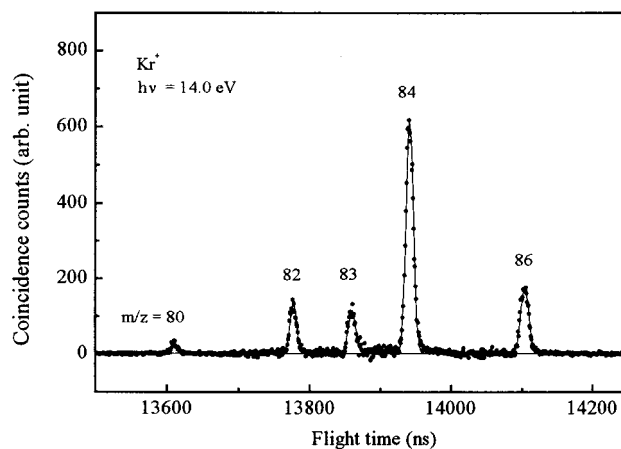
**2.1. Experiment.** Coincidence measurements were performed with a molecular-beam/threshold-photoelectron-photoion-coincidence (MB/TPEPICO) apparatus, described in detail elsewhere.<sup>25</sup> In brief, synchrotron radiation from the 1.5 GeV electron storage ring at the Synchrotron Radiation Research Center (SRRC) in Taiwan is dispersed with a 1 m Seya-Namioka monochromator having a 1200 lines/mm grating to cover an energy range of 11–12 eV. A resolution of  $\sim 0.02$  eV and a photon flux of  $> 10^9$  photons  $\text{s}^{-1}$  were achieved when both entrance and exit slits were set at 0.1 mm. Absolute photon energies were calibrated to within  $\pm 0.003$  eV on measurement of well-characterized Rydberg peaks in the TPES spectra of Ar and Kr.

In the ionization chamber, the monochromatic vacuum ultraviolet (VUV) radiation intersected perpendicularly with a molecular beam, formed on expanding a  $\text{CH}_3\text{CFCl}_2/\text{He}$  mixture through a nozzle and two skimmers. The total stagnation pressure was  $\sim 290$  Torr, with a seed ratio of  $\sim 10\%$ . A threshold photoelectron spectrometer and a time-of-flight (TOF) mass spectrometer of Wiley–McLaren type were mounted in opposite directions to measure threshold electrons and ions, respectively; dual microchannel plates served as detectors in both spectrometers. Axes of both spectrometers were perpendicular to those of molecular and photon beams. To improve instrumental resolutions and efficiency of ion collection, a pulsed field of  $40 \text{ V cm}^{-1}$  with a duration of  $30 \mu\text{s}$  was applied to extract ions whereas electrons were accelerated with a dc field of  $1.64 \text{ V cm}^{-1}$  before detection.

Signals of detected threshold electrons and ions were fed into a time-to-digital converter (TDC) to record flight durations of ions that were detected within  $30 \mu\text{s}$  of each cycle triggered with a threshold electron. TOF distributions of ions in coincidence with triggering threshold electrons were obtained after  $(5\text{--}15) \times 10^5$  cycles depending on coincident ion counts. Because synchrotron radiation was operated at a repetition rate of  $\sim 500$  MHz, each measured coincidence spectrum has a large contribution from uncorrelated ions piling up in the interaction region. To eliminate such a contribution, a second coincidence spectrum triggered with a signal generated randomly relative to the preceding threshold electron signal was accumulated following each threshold electron-triggered coincidence cycle. Subtraction of the randomly generated coincidence spectrum from the electron-triggered coincidence spectrum yielded a true coincidence spectrum. All data acquisition was controlled with a computer via CAMAC interface, and output from the TDC converter was transferred to the computer for further processing.

$\text{CH}_3\text{CFCl}_2$  (Darkin Corp., 99.9%) was degassed with several freeze–pump–thaw cycles before use. Helium ( $> 99.999\%$ ), krypton ( $> 99.99\%$ ), and argon ( $> 99.99\%$ ) were used without further purification.

**2.2. Theoretical Calculations.** We calculated ionization energies of  $\text{CH}_3\text{CFCl}_2$  and  $\text{CH}_3\text{CFCl}$  and enthalpy changes of isodesmic reactions related to formation of these two molecules with the GAUSSIAN 98 program.<sup>26</sup> The equilibrium structures were fully optimized using Schegel's analytical gradient method<sup>27</sup> with a B3LYP density functional<sup>28,29</sup> which is a combination of Slater exchange, Becke exchange,<sup>30</sup> and Lee–Yang–Parr (LYP) correlation functionals.<sup>31</sup> Standard basis sets 6-31+G(d) and aug-cc-pVTZ,<sup>32</sup> a Dunning's correlation-consistent polarized valence triple- $\zeta$  basis set augmented with s, p, d, and



**Figure 1.** Coincidence TOF spectrum of Kr excited at photon energy 14.0 eV.

f diffuse functions, were used in all structural optimization and calculations of vibrational frequencies.

The G3B3 theory<sup>33</sup> is based on B3LYP/6-31G(d) geometry and zero-point energies scaled by 0.96. All other steps remain the same as G3 theory except values of higher-level correction parameters,  $A = 6.760$  mhartree,  $B = 3.233$  mhartree,  $C = 6.786$  mhartree, and  $D = 1.269$  mhartree. This G3B3 theory was used to calculate the ionization energy and to predict the heat of formation of  $\text{CH}_3\text{CFCl}_2$  and  $\text{CH}_3\text{CFCl}$  via isodesmic reactions in which the types and number of bonds are the same for products and reactants.

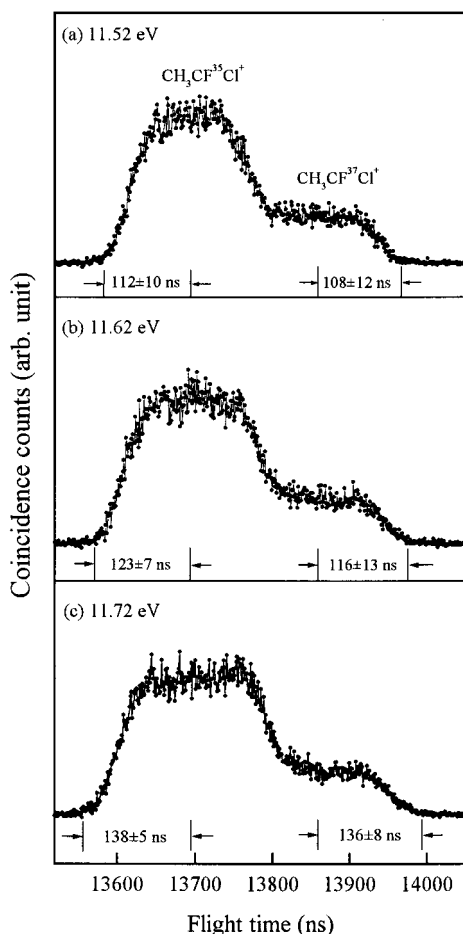
## 3. Results and Discussion

**3.1. Coincidence TOF Spectra for  $\text{CH}_3\text{CFCl}^+$ .** Coincidence TOF spectra for helium, argon, and krypton excited at photon energies near their ionization thresholds were recorded to calibrate ion flight times and to verify the effect of kinetic energy release on TOF bandwidths upon dissociation. A coincidence spectrum of krypton excited at 14.0 eV is shown in Figure 1; five TOF peaks correspond to krypton isotopes at  $m/z = 80$ , 82, 83, 84, and 86. Each peak in this coincidence spectrum was fitted to a Gaussian profile with a full width at half-maximum (fwhm) of  $\sim 12$  ns, indicating a small transverse velocity of the skimmed molecular beam along the detection axis. Observed flight times of 3218 ( $\text{He}^+$ ), 9689 ( $\text{Ar}^+$ ), 13610 ( $^{80}\text{Kr}^+$ ), 13776 ( $^{82}\text{Kr}^+$ ), 13859 ( $^{83}\text{Kr}^+$ ), 13941 ( $^{84}\text{Kr}^+$ ), and 14103 ns ( $^{86}\text{Kr}^+$ ) were fitted to an equation

$$T_0 = 1496.5 (m/z)^{1/2} + 225.3 \quad (1)$$

in which  $T_0$  is ion flight time in ns. Flight times calculated with eq 1 agree within 1 ns of experimental values.

Coincidence TOF spectra of  $\text{CH}_3\text{CFCl}_2$  excited above appearance onset of  $\text{CH}_3\text{CFCl}^+$  at 11.33 eV were recorded at photon energies of 11.42, 11.52, 11.62, 11.72, and 11.82 eV with 1 ns resolution and corrected for false coincidence background. Corrected coincidence spectra in the region of having ion coincident counts obtained after excitation of  $\text{CH}_3\text{CFCl}_2$  at photon energies of 11.52, 11.62, and 11.72 eV are shown in Figure 2a–c; all spectra are normalized to an observed maximal intensity. All spectra consist of two partially overlapping, symmetrically broadened, and nearly flat-topped TOF peaks. The intensity ratios of these two TOF peaks are near the  $^{35}\text{Cl}/^{37}\text{Cl}$  isotopic ratio of  $\sim 3$ , and their respective fwhm are greater than 150 ns. We assigned these two TOF peaks as  $\text{CH}_3$ -



**Figure 2.** Coincidence TOF spectra of  $\text{CH}_3\text{CFCl}_2$  excited at photon energies (a) 11.52, (b) 11.62, and (c) 11.72 eV.

$\text{CF}^{35}\text{Cl}^+$  ( $m/z = 81$ ) and  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$  ( $m/z = 83$ ) on the basis of observed flight times and their intensity ratio.

### 3.2. Thermochemistry of the $\text{CH}_3\text{CFCl}^+ + \text{Cl}$ Channel.

The maximal releases of kinetic energy with respect to the center of mass ( $\text{KE}_{\text{CM}}$ ) of  $\text{CH}_3\text{CFCl}^+$  were derived from forward half-widths of  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  TOF peaks and backward half-widths of  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$  TOF peaks in the coincidence spectra measured at five photon energies. The full width of a TOF peak corresponds to the turnaround period of the most energetic ion moving away from the TOF ion detector; the analysis of half-widths is required due to overlapping  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  and  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$  TOF peaks. We calculated  $T_0 = 13694$  and  $13859$  ns from eq 1 for  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  and  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$ ; half-widths of  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  and  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$  TOF peaks subsequently yield  $\text{KE}_{\text{CM}}$  for  $\text{CH}_3\text{CFCl}^+$  and total kinetic energy releases ( $\text{KE}_{\text{T}}$ ) for a channel of formation of  $\text{CH}_3\text{CFCl}^+$  according to the following equations:

$$\text{KE}_{\text{CM}}(\text{CH}_3\text{CFCl}^+) = [1/(2M_{\text{CH}_3\text{CFCl}})](\Delta tqV)^2 \quad (2)$$

in which  $M_{\text{CH}_3\text{CFCl}} = 81$  or  $83$  g mol $^{-1}$ ,  $\Delta t$  is the half-width of the corresponding TOF peak,  $q$  is the charge of  $\text{CH}_3\text{CFCl}^+$ , and  $V = 41.64$  V cm $^{-1}$  is the electric field, and

$$\text{KE}_{\text{T}} = \text{KE}_{\text{CM}}(\text{CH}_3\text{CFCl}^+)(M_{\text{CH}_3\text{CFCl}_2}/M_{\text{Cl}}) \quad (3)$$

in which  $M_{\text{CH}_3\text{CFCl}_2} = 116$  or  $118$  or  $120$ , and  $M_{\text{Cl}} = 35$  or  $37$  g mol $^{-1}$ .

We assume that bond energies and rates of bond breaking of  $\text{C}-^{35}\text{Cl}$  and  $\text{C}-^{37}\text{Cl}$  are the same as those in the calculations

of  $\text{KE}_{\text{T}}$ , although formation of  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  is due to loss of  $^{35}\text{Cl}$  from  $\text{CH}_3\text{CF}^{35}\text{Cl}^{35}\text{Cl}^+$  and loss of  $^{37}\text{Cl}$  from  $\text{CH}_3\text{CF}^{35}\text{Cl}^{37}\text{Cl}^+$ . Because  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  TOF peaks originating from such an isotopic effect are unresolved,  $\text{KE}_{\text{T}}$  for losses of  $^{35}\text{Cl}$  and  $^{37}\text{Cl}$  at one-photon energy were calculated from eq 3, and their average value was adopted for formation of  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$ . The uncertainty of this treatment, less than  $0.02$  eV, is smaller than that of  $0.09$  eV from determination of half-widths of  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  TOF peaks. The  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$  TOF peaks resulting from loss of  $^{35}\text{Cl}$  from  $\text{CH}_3\text{CF}^{35}\text{Cl}^{37}\text{Cl}^+$  and loss of  $^{37}\text{Cl}$  from  $\text{CH}_3\text{CF}^{37}\text{Cl}^{37}\text{Cl}^+$  were analyzed in a similar way. The determined half-widths for TOF peaks of  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  and  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$ , and corresponding  $\text{KE}_{\text{T}}$  derived at five photon energies, are listed in Table 1.

AE values for the  $\text{CH}_3\text{CFCl}^+ + \text{Cl}$  channel at five photon energies were determined by subtracting corresponding  $\text{KE}_{\text{T}}$  from photon energy; they are also listed in Table 1. An average value of  $11.10 \pm 0.09$  eV is derived for the  $\text{CH}_3\text{CF}^{35}\text{Cl}^+ + ^{35}\text{Cl}/^{37}\text{Cl}$  channel, whereas a value of  $11.13 \pm 0.16$  eV is derived for the  $\text{CH}_3\text{CF}^{37}\text{Cl}^+ + ^{35}\text{Cl}/^{37}\text{Cl}$  channel. Listed errors are mainly due to uncertainties on determination of half-widths. The values of AE for both channels agree satisfactorily, except that the determined AE from the latter channel suffers from large errors because of small  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$  signals. Thus, we adopted  $\text{AE} = 11.10 \pm 0.09$  eV. The AE is the thermochemical threshold of the  $\text{CH}_3\text{CFCl}^+ + \text{Cl}$  channel if no reverse barrier exists for this process; accordingly  $\text{KE}_{\text{T}} = 0.23$  eV at the appearance onset of  $\text{CH}_3\text{CFCl}^+$  was derived on subtraction of the determined AE from the appearance onset of  $11.33$  eV of  $\text{CH}_3\text{CFCl}^+$  determined in a PIMS experiment.<sup>17</sup>

The determined AE relates the heats of formation of  $\text{CH}_3\text{CFCl}^+$ ,  $\text{Cl}$ , and  $\text{CH}_3\text{CFCl}_2$ :

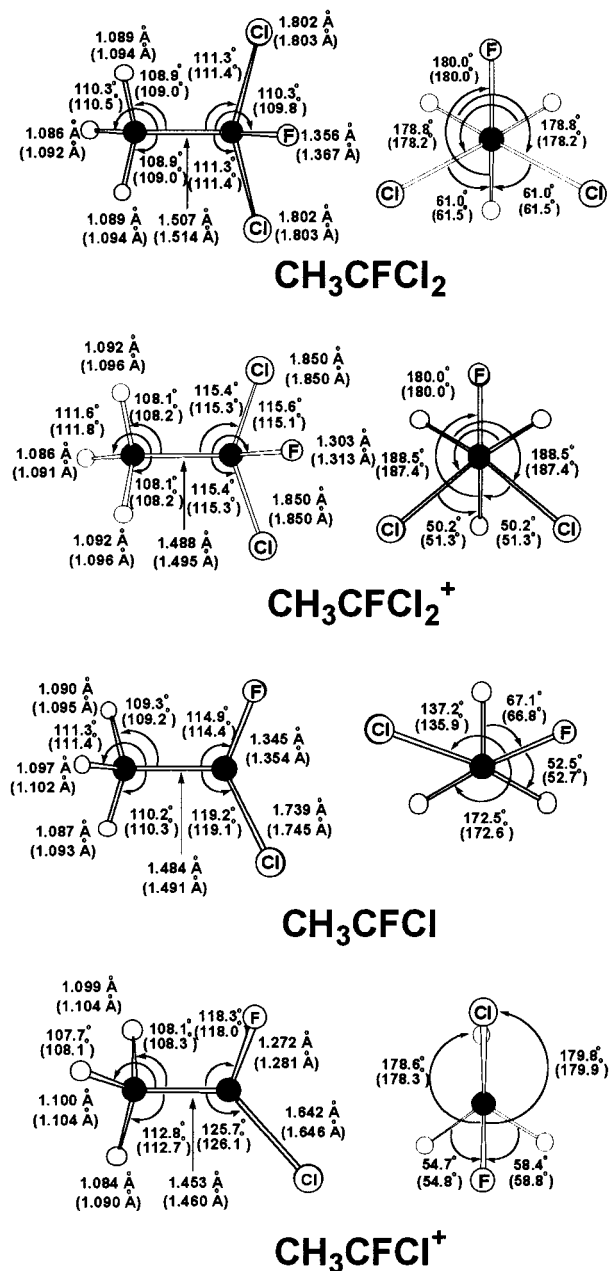
$$\text{AE} = \Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+) + \Delta_f H_0^\circ(\text{Cl}) - \Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2) - E_{\text{int}} \quad (4)$$

in which  $\Delta_f H_0^\circ$  is the heat of formation at  $0$  K and  $E_{\text{int}}$  is the average internal energy of  $\text{CH}_3\text{CFCl}_2$ . Thus, to derive  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+)$  from eq 4 with the determined AE value,  $\Delta_f H_0^\circ(\text{Cl})$  and  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2)$  are required.  $E_{\text{int}}$  is neglected because  $\text{CH}_3\text{CFCl}_2$  is cooled in the molecular beam. A literature value for  $\Delta_f H_0^\circ(\text{Cl})$  is  $28.6$  kcal mol $^{-1}$ .<sup>34</sup> Although an experimental value for  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2)$  is unavailable, we estimate it by interpolation with a nonlinear fit of literature values of  $\Delta_f H_{298}^\circ$  for  $\text{CH}_3\text{CCl}_3$ ,  $\text{CH}_3\text{CF}_2\text{Cl}$ , and  $\text{CH}_3\text{CF}_3$ . Values of  $\Delta_f H_{298}^\circ$  for  $\text{CH}_3\text{CCl}_3$ ,  $\text{CH}_3\text{CF}_2\text{Cl}$ , and  $\text{CH}_3\text{CF}_3$  are  $-34.54$ ,  $-128.09$ , and  $-178.86$  kcal mol $^{-1}$ , respectively,<sup>34</sup> an estimated  $\Delta_f H_{298}^\circ(\text{CH}_3\text{CFCl}_2)$  is  $-80.1$  kcal mol $^{-1}$ . We tested this method for analogous compounds in two series,  $\text{CHF}_x\text{Cl}_{3-x}$  and  $\text{CF}_x\text{Cl}_{3-x}$  ( $x = 0-3$ ),<sup>35</sup> and found the variation between literature and fitted values to be less than  $0.16$  kcal mol $^{-1}$ . Hence,  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2) = -78.4$  kcal mol $^{-1}$  was derived after correction of  $\Delta_f H_{298}^\circ(\text{CH}_3\text{CFCl}_2)$  using experimental vibrational frequencies of  $3028$ ,  $3012$ ,  $2954$ ,  $1442$ ,  $1442$ ,  $1385$ ,  $1159$ ,  $1119$ ,  $1092$ ,  $926$ ,  $750$ ,  $591$ ,  $432$ ,  $398$ ,  $380$ ,  $291$ ,  $262$ , and  $248$  cm $^{-1}$  for  $\text{CH}_3\text{CFCl}_2$ .<sup>36</sup> Accordingly,  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+) = 149.0 \pm 2.1$  kcal mol $^{-1}$  was derived from eq 4.

**3.3. Theoretical Predictions.** We employed theoretical calculations with the G3B3 method to derive  $\Delta_f H_0^\circ$  for  $\text{CH}_3\text{CFCl}_2$ ,  $\text{CH}_3\text{CFCl}$ , and their corresponding cations. Structures of  $\text{CH}_3\text{CFCl}_2$ ,  $\text{CH}_3\text{CFCl}_2^+$ ,  $\text{CH}_3\text{CFCl}$ , and  $\text{CH}_3\text{CFCl}^+$  optimized with B3LYP method are shown in Figure 3. Results using the aug-cc-pVTZ basis set are listed, whereas those using the 6-31+G(d) basis set are listed in parentheses for comparison.

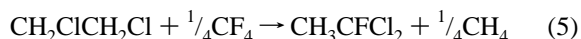
**TABLE 1: Half-Widths ( $\Delta t_{\text{half-width}}$ ) of  $\text{CH}_3\text{CF}^{35}\text{Cl}^+$  and  $\text{CH}_3\text{CF}^{37}\text{Cl}^+$  TOF Peaks, Total Kinetic Energy Release ( $\text{KE}_T$ ), and Appearance Energy (AE) Measured at Five Photon Energies (PE) near the Appearance Onset of  $\text{CH}_3\text{CFCl}^+$** 

PE/eV	$\text{CH}_3\text{CF}^{35}\text{Cl}^+$			$\text{CH}_3\text{CF}^{37}\text{Cl}^+$		
	$\Delta t_{\text{half-width}}/\text{ns}$	$\text{KE}_T/\text{eV}$	AE/eV	$\Delta t_{\text{half-width}}/\text{ns}$	$\text{KE}_T/\text{eV}$	AE/eV
11.42	99 ± 9	0.33 ± 0.05	11.09 ± 0.05	102 ± 9	0.35 ± 0.06	11.07 ± 0.06
11.52	112 ± 10	0.42 ± 0.07	11.10 ± 0.07	108 ± 12	0.39 ± 0.09	11.13 ± 0.09
11.62	123 ± 7	0.50 ± 0.07	11.12 ± 0.07	116 ± 13	0.44 ± 0.11	11.18 ± 0.11
11.72	138 ± 5	0.63 ± 0.04	11.09 ± 0.04	136 ± 8	0.61 ± 0.07	11.11 ± 0.07
11.82	144 ± 6	0.70 ± 0.05	11.12 ± 0.05	142 ± 7	0.67 ± 0.06	11.15 ± 0.06
			Av. 11.10 ± 0.09			Av. 11.13 ± 0.16

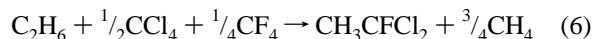
**Figure 3.** Structures of  $\text{CH}_3\text{CFCl}_2$ ,  $\text{CH}_3\text{CFCl}_2^+$ ,  $\text{CH}_3\text{CFCl}$ , and  $\text{CH}_3\text{CFCl}^+$  optimized with B3LYP/aug-cc-pVTZ method; results using 6-31+G(d) basis sets are listed in parentheses.

For enthalpy changes ( $\Delta H^\circ$ ) of three isodesmic reactions and ionization energies (IE) of  $\text{CH}_3\text{CFCl}_2$  and  $\text{CH}_3\text{CFCl}$ , Table 2 lists resultant  $\Delta H^\circ$  or IE, literature values<sup>34,35</sup> of  $\Delta_f H_0^\circ$  for species relevant to calculations, and derived  $\Delta_f H_0^\circ$  for species of interest.

In Table 2, the enthalpy change  $\Delta H^\circ$  for the isodesmic reaction

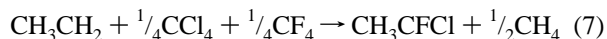


was calculated to be 2.0 kcal mol<sup>-1</sup>, consequently  $\Delta_f H_0^\circ$  ( $\text{CH}_3\text{CFCl}_2$ ) = -78.4 kcal mol<sup>-1</sup> was derived using literature values<sup>34,35</sup> of  $\Delta_f H_0^\circ$  for  $\text{CH}_2\text{ClCH}_2\text{Cl}$ ,  $\text{CF}_4$ , and  $\text{CH}_4$ . Similarly, we derive  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2) = -77.4$  kcal mol<sup>-1</sup> from an isodesmic reaction in which calculated  $\Delta H^\circ$  is -6.4 kcal mol<sup>-1</sup>.



The average value of -77.9 kcal mol<sup>-1</sup> was used for  $\Delta_f H_0^\circ$  ( $\text{CH}_3\text{CFCl}_2$ ). This value varies by only 1 kcal mol<sup>-1</sup> from a value of -76.93 kcal mol<sup>-1</sup> derived with the ab initio BAC-MP4 method<sup>37</sup> and 0.5 kcal mol<sup>-1</sup> from a value of -78.4 kcal mol<sup>-1</sup> estimated from interpolation of  $\Delta_f H_{298}^\circ$  of  $\text{CH}_3\text{CF}_x\text{Cl}_{3-x}$  ( $x = 0-3$ ), as described above. In addition, IE of  $\text{CH}_3\text{CFCl}_2$  was calculated to be 255.1 kcal mol<sup>-1</sup>, and accordingly  $\Delta_f H_0^\circ$  ( $\text{CH}_3\text{CFCl}_2^+$ ) = 177.2 kcal mol<sup>-1</sup> was derived.

Also listed in Table 2, calculation on the isodesmic reaction



yields  $\Delta H^\circ = -0.4$  kcal mol<sup>-1</sup>. Combining this  $\Delta H^\circ$  with literature values<sup>34,35</sup> of  $\Delta_f H_0^\circ$  for  $\text{CH}_3\text{CH}_2$ ,  $\text{CCl}_4$ ,  $\text{CF}_4$ , and  $\text{CH}_4$ , we derive  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}) = -22.3$  kcal mol<sup>-1</sup>. An IE of 173.8 kcal mol<sup>-1</sup> for  $\text{CH}_3\text{CFCl}$  was also calculated, then  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+)$  was derived to be 151.5 kcal mol<sup>-1</sup>. This value is near the experimental value of  $149.0 \pm 2.1$  kcal mol<sup>-1</sup> determined in this work.

Curtiss et al. made test calculations on dissociation energies, IE, and electron affinities using G3, G3(MP2), and G3(MP3) methods on 376 energies, and they found that the average absolute deviations of calculated values from experimental values are 1.07, 1.31, and 1.27 kcal mol<sup>-1</sup>, respectively.<sup>38</sup> The uncertainty associated with calculations on isodesmic reactions is typically  $\pm 1$  kcal mol<sup>-1</sup>. Thus, our experimental value of  $149.0 \pm 2.1$  kcal mol<sup>-1</sup> for  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+)$  agrees with the theoretical value of 151.5 kcal mol<sup>-1</sup>.

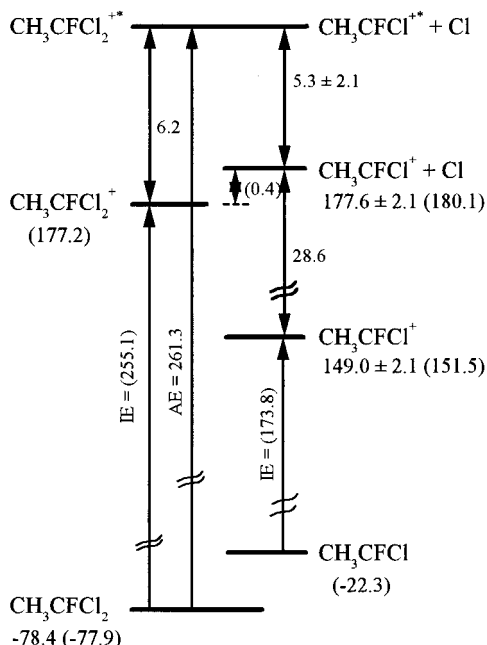
**3.4. Dissociative Photoionization of  $\text{CH}_3\text{CFCl}_2$  To Form  $\text{CH}_3\text{CFCl}^+ + \text{Cl}$ .** Observation of only  $\text{CH}_3\text{CFCl}^+$  with no parent ion  $\text{CH}_3\text{CFCl}_2^+$  in the photoionization energy region 11.42–11.82 eV is consistent with the PIMS experiment. The broadening of  $\text{CH}_3\text{CFCl}^+$  TOF peaks in the coincidence spectra is due to release of kinetic energy upon dissociative photoionization of  $\text{CH}_3\text{CFCl}_2$ , as explained above. In an ideal case, a single value for released kinetic energy and an isotropic distribution would give rise to a square TOF profile when the recoiling velocity of fragments is greater than that of a parent ion. A Gaussian shape of TOF peaks would be observed when the distribution of released kinetic energy is nearly Maxwellian upon dissociation of an ion. Thus, observed nearly flat-top and broadened TOF profiles imply that the distribution of kinetic energy of  $\text{CH}_3\text{CFCl}^+$  is narrow, characteristic of a repulsive potential energy surface for the dissociation of  $\text{CH}_3\text{CFCl}_2^+$  into  $\text{CH}_3\text{CFCl}^+$ .



**TABLE 2: Ionization Energies (IE) of CH<sub>3</sub>CFCl<sub>2</sub> and CH<sub>3</sub>CFCl and Heats of Formation ( $\Delta_f H_0^\circ$ ) for CH<sub>3</sub>CFCl<sub>2</sub>, CH<sub>3</sub>CFCl<sub>2</sub><sup>+</sup>, CH<sub>3</sub>CFCl, and CH<sub>3</sub>CFCl<sup>+</sup> Derived from Calculations of Isodesmic Reactions Based on G3B3 Theory, and Literature Values of  $\Delta_f H_0^\circ$  for Relevant Species**

equations	$\Delta H^\circ$ or IE /kcal mol <sup>-1</sup>	literature $\Delta_f H_0^\circ$ <sup>a</sup>		derived $\Delta_f H_0^\circ$ <sup>a</sup>	
		species	/kcal mol <sup>-1</sup>	species	/kcal mol <sup>-1</sup>
1 CH <sub>2</sub> ClCH <sub>2</sub> Cl + 1/4 CF <sub>4</sub> → CH <sub>3</sub> CFCl <sub>2</sub> + 1/4 CH <sub>4</sub>	$\Delta H^\circ = 2.0$	CH <sub>2</sub> ClCH <sub>2</sub> Cl CF <sub>4</sub> CH <sub>4</sub>	-29 -221.6 <sup>b</sup> -16.0 <sup>b</sup>	CH <sub>3</sub> CFCl <sub>2</sub>	-78.4
2 CH <sub>3</sub> CH <sub>3</sub> + 1/2 CCl <sub>4</sub> + 1/4 CF <sub>4</sub> → CH <sub>3</sub> CFCl <sub>2</sub> + 3/4 CH <sub>4</sub>	$\Delta H^\circ = -6.4$	CH <sub>3</sub> CH <sub>3</sub> CCl <sub>4</sub>	-16.4 -22.4 <sup>b</sup>	CH <sub>3</sub> CFCl <sub>2</sub>	-77.4
3 CH <sub>3</sub> CH <sub>2</sub> + 1/4 CCl <sub>4</sub> + 1/4 CF <sub>4</sub> → CH <sub>3</sub> CClF + 1/2 CH <sub>4</sub>	$\Delta H^\circ = -0.4$	CH <sub>3</sub> CH <sub>2</sub>	31.1	CH <sub>3</sub> CFCl	-22.3
4 CH <sub>3</sub> CFCl <sub>2</sub> → CH <sub>3</sub> CFCl <sub>2</sub> <sup>+</sup> + e <sup>-</sup>	IE = 255.1	CH <sub>3</sub> CFCl <sub>2</sub>	-77.9 <sup>c</sup>	CH <sub>3</sub> CFCl <sub>2</sub> <sup>+</sup>	177.2
5 CH <sub>3</sub> CFCl → CH <sub>3</sub> CFCl <sup>+</sup> + e <sup>-</sup>	IE = 173.8	CH <sub>3</sub> CFCl	-22.3	CH <sub>3</sub> CFCl <sup>+</sup>	151.5

<sup>a</sup> Reference 35, unless stated otherwise. <sup>b</sup> Reference 34. <sup>c</sup> An average value of  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2)$  from reactions 1 and 2 was used.



**Figure 4.** Schematic energy diagram for dissociative photoionization of CH<sub>3</sub>CFCl<sub>2</sub> to form CH<sub>3</sub>CFCl<sup>+</sup> + Cl based on experimental and calculated results; the values are in kcal mol<sup>-1</sup>, and calculated  $\Delta_f H_0^\circ$  or IE are listed in parentheses.

Figure 4 shows an energy diagram for dissociative photoionization of CH<sub>3</sub>CFCl<sub>2</sub> to form CH<sub>3</sub>CFCl<sup>+</sup> + Cl based on experimental and calculated results; calculated  $\Delta_f H_0^\circ$  or IE are listed in parentheses. The experimentally determined thermochemical threshold for CH<sub>3</sub>CFCl<sup>+</sup> + Cl channel is  $11.10 \pm 0.09$  eV ( $256.0 \pm 2.1$  kcal mol<sup>-1</sup>), and accordingly the energy of this channel is  $177.6$  kcal mol<sup>-1</sup> using  $\Delta_f H_0^\circ(\text{Cl}) = 28.6$  kcal mol<sup>-1</sup> and  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2) = -78.4$  kcal mol<sup>-1</sup>. The theoretically predicted  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2^+)$  is  $177.2$  kcal mol<sup>-1</sup>. The variation of  $0.4$  kcal mol<sup>-1</sup> between energies of CH<sub>3</sub>CFCl<sub>2</sub><sup>+</sup> and CH<sub>3</sub>CFCl<sup>+</sup> + Cl is smaller than uncertainties in experiments and theoretical calculations. Hence, we conclude that the dissociation process



is nearly thermoneutral.

Alternatively, the energy of CH<sub>3</sub>CFCl<sup>+</sup> + Cl can be estimated from ionization of CH<sub>3</sub>CFCl. Combining predicted values of  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}) = -22.3$  kcal mol<sup>-1</sup>, IE(CH<sub>3</sub>CFCl) =  $173.8$  kcal mol<sup>-1</sup>, and  $\Delta_f H_0^\circ(\text{Cl}) = 28.6$  kcal mol<sup>-1</sup>, we obtain an energy of  $180.1$  kcal mol<sup>-1</sup> for CH<sub>3</sub>CFCl<sup>+</sup> + Cl. This value is greater than the value  $177.6$  kcal mol<sup>-1</sup> by  $2.5$  kcal mol<sup>-1</sup>, but

still within possible error limits of experiments and calculations. Both experimental and predicted values of the energy of CH<sub>3</sub>CFCl<sup>+</sup> + Cl are slightly greater than  $177.2$  kcal mol<sup>-1</sup> predicted for  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}_2^+)$ , indicating that reaction 8 is perhaps slightly endothermic.

An appearance onset of  $11.33$  eV ( $261.3$  kcal mol<sup>-1</sup>) for CH<sub>3</sub>CFCl<sup>+</sup>, about  $6.2$  kcal mol<sup>-1</sup> above the predicted IE of  $255.1$  kcal mol<sup>-1</sup> for CH<sub>3</sub>CFCl<sub>2</sub>, was determined in the PIMS experiments. According to calculations, geometries of CH<sub>3</sub>CFCl<sub>2</sub> and CH<sub>3</sub>CFCl<sub>2</sub><sup>+</sup> differ significantly; C–Cl bond lengths of CH<sub>3</sub>CFCl<sub>2</sub> and CH<sub>3</sub>CFCl<sub>2</sub><sup>+</sup> are  $1.80$  and  $1.85$  Å, and Cl–C–Cl angles are  $122.0^\circ$  and  $100.4^\circ$ , respectively, as shown in Figure 3. Detection of a KE<sub>T</sub> of  $0.23$  eV ( $5.3$  kcal mol<sup>-1</sup>) at the appearance onset of CH<sub>3</sub>CFCl<sup>+</sup> and observation of only CH<sub>3</sub>CFCl<sup>+</sup> with no CH<sub>3</sub>CFCl<sub>2</sub><sup>+</sup> also support that reaction 8 is nearly thermoneutral or slightly endothermic, and that CH<sub>3</sub>CFCl<sub>2</sub><sup>+</sup>, produced immediately after CH<sub>3</sub>CFCl<sub>2</sub> absorbs a VUV photon, undergoes dissociation to form CH<sub>3</sub>CFCl<sup>+</sup>.

Similar phenomena were observed for SF<sub>6</sub><sup>+</sup> and CF<sub>4</sub><sup>+</sup> by Creasey et al. In these cases, the parent ions are virtually undetectable because their ground states are repulsive and lie above the lowest dissociation channel.<sup>39,40</sup> Creasey et al. found that, for SF<sub>6</sub><sup>+</sup>, excess energies above the dissociation threshold are partitioned into translational energy and internal energy of fragments via a direct and rapid dissociation path; for this reason statistical methods cannot be applied to model the energy distribution. Our observed nearly flat-top and broadened TOF profiles are consistent with their observation.

#### 4. Conclusions

With a threshold photoelectron photoion coincidence technique in pulsed mode, we investigated the formation of CH<sub>3</sub>CFCl<sup>+</sup> upon photoionization of CH<sub>3</sub>CFCl<sub>2</sub> at photon energies near the appearance onset of CH<sub>3</sub>CFCl<sup>+</sup>. After correction of maximal kinetic energy releases determined from half-widths of CH<sub>3</sub>CF<sup>35</sup>Cl<sup>+</sup> TOF peaks in coincidence spectra, we derive a thermochemical threshold at  $11.10 \pm 0.09$  eV for the CH<sub>3</sub>CFCl<sup>+</sup> + Cl channel, a maximal kinetic energy release of  $0.23$  eV at the appearance onset of CH<sub>3</sub>CFCl<sup>+</sup>, and  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+) = 149.0 \pm 2.1$  kcal mol<sup>-1</sup>. Enthalpy changes of three isodesmic reactions related to formation of CH<sub>3</sub>CFCl<sub>2</sub> and CH<sub>3</sub>CFCl, and ionization energies of CH<sub>3</sub>CFCl<sub>2</sub> and CH<sub>3</sub>CFCl, were also computed to derive heats of formation of CH<sub>3</sub>CFCl<sub>2</sub>, CH<sub>3</sub>CFCl, and their corresponding cations; observed  $\Delta_f H_0^\circ(\text{CH}_3\text{CFCl}^+)$  agrees satisfactorily with calculation. An energy diagram is constructed to rationalize the dissociative photoionization of CH<sub>3</sub>CFCl<sub>2</sub> near the threshold region. Dissociation of CH<sub>3</sub>CFCl<sub>2</sub><sup>+</sup> into CH<sub>3</sub>CFCl<sup>+</sup> + Cl is nearly thermoneutral, perhaps slightly endothermic; the excess energy in CH<sub>3</sub>CFCl<sub>2</sub><sup>+</sup> produced im-

mediately after photoionization facilitates dissociation, consistent with lack of observation of  $\text{CH}_3\text{CFCl}_2^+$ .

**Acknowledgment.** We thank the Synchrotron Radiation Research Center and the National Science Council of Taiwan for financial support of this work (contract no. NSC89-2113-M-213-008) and the National Center for High-performance Computing for providing computing time for theoretical calculations.

## References and Notes

- (1) Schaufli, S. M.; Pollock, W. H.; Atlas, E. L.; Heidt, L. E.; Daniel, J. S. *Geophys. Res. Lett.* **1995**, *22*, 819.
- (2) Gehring, D. G.; Barsotti, D. J.; Gibbon, H. E. *J. Chromatogr. Sci.* **1992**, *30*, 301.
- (3) Talukdar, R.; Mellouki, A.; Gierczak, T.; Burkholder, J. B.; McKeen, S. A.; Ravishankara, A. R. *J. Phys. Chem.* **1991**, *95*, 5815.
- (4) Fahr, A.; Braun, W.; Kurylo, M. J. *J. Geophys. Res.* **1993**, *98*, 20467.
- (5) Gillotay, D.; Simon, P. C.; Brasseur, G. *Planet. Space Sci.* **1989**, *37*, 105.
- (6) Melchior, A.; Bar, I.; Rosenwaks, S. *J. Chem. Phys.* **1997**, *107*, 8476.
- (7) Melchior, A.; Chen, X.; Bar, I.; Rosenwaks, S. *Chem. Phys. Lett.* **1999**, *315*, 421.
- (8) Huder, K.; DeMore, W. B. *Geophys. Res. Lett.* **1993**, *20*, 1575.
- (9) Liu, R. Z.; Huie, R. E.; Kurylo, M. J. *J. Phys. Chem.* **1990**, *94*, 3247.
- (10) Zhang, Z. Y.; Huie, R. E.; Kurylo, M. J. *J. Phys. Chem.* **1992**, *96*, 1533.
- (11) Warren, R.; Gierczak, T.; Ravishankara, A. R. *Chem. Phys. Lett.* **1991**, *183*, 403.
- (12) Tuazon, E. C.; Atkinson, R.; Corchnoy, S. B. *Int. J. Chem. Kinet.* **1992**, *24*, 639.
- (13) Sawerysyn, J. P.; Talhaoui, A.; Meriaux, B.; Devolder, P. *Chem. Phys. Lett.* **1992**, *198*, 197.
- (14) Talhaoui, A.; Louis, F.; Devolder, P.; Meriaux, B.; Sawerysyn, J.-P.; Rayez, M.-T.; Rayez, J.-C. *J. Phys. Chem.* **1996**, *100*, 13531.
- (15) Tuazon, E. C.; Atkinson, R. *J. Atmos. Chem.* **1993**, *17*, 179.
- (16) Warren, R. F.; Ravishankara, A. R. *Int. J. Chem. Kinet.* **1993**, *25*, 833.
- (17) Chiang, S.-Y.; Wang, T.-T.; Yu, J.-S. K.; Yu, C.-h. *Chem. Phys. Lett.* **2000**, *329*, 185.
- (18) Tsai, B. P.; Baer, T.; Werner, A. S.; Lin, S. F. *J. Phys. Chem.* **1975**, *79*, 570.
- (19) Shiromaru, H.; Achiba, Y.; Kimura, K.; Lee, Y. T. *J. Phys. Chem.* **1987**, *91*, 17.
- (20) Baer, T. *Adv. Chem. Phys.* **1986**, *64*, 111.
- (21) Mayer, P. M.; Baer, T. *Chem. Phys. Lett.* **1996**, *261*, 155.
- (22) Evans, M.; Ng, C. Y.; Hsu, C.-W.; Heimann, P. *J. Chem. Phys.* **1997**, *106*, 978.
- (23) Cheung, Y. S.; Chen, Y.-J.; Ng, C. Y.; Chiu, S. W.; Li, W.-K. *J. Am. Chem. Soc.* **1995**, *117*, 9725.
- (24) Curtiss, L. A.; Raghavachari, K.; Pople, J. A. *J. Chem. Phys.* **1993**, *98*, 1293.
- (25) Chiang, S.-Y.; Ma, C.-I. *J. Phys. Chem. A* **2000**, *104*, 1991.
- (26) Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseria, G. E.; Robb, M. A.; Cheeseman, J. R.; Zakrzewski, V. G.; Montgomery, J. A. Jr.; Stratmann, R. E.; Burant, J. C.; Dapprich, S.; Millam, J. M.; Daniels, A. D.; Kudin, K. N.; Strain, M. C.; Farkas, O.; Tomasi, J.; Barone, V.; Cossi, M.; Cammi, R.; Mennucci, B.; Pomelli, C.; Adamo, C.; Clifford, S.; Ochterski, J.; Petersson, G. A.; Ayala, P. Y.; Cui, Q.; Morokuma, K.; Malick, D. K.; Rabuck, A. D.; Raghavachari, K.; Foresman, J. B.; Cioslowski, J.; Ortiz, J. V.; Baboul, A. G.; Stefanov, B. B.; Liu, G.; Liashenko, A.; Piskorz, P.; Komaromi, I.; Gomperts, R.; Martin, R. L.; Fox, D. J.; Keith, T.; Al-Laham, M. A.; Peng, C. Y.; Nanayakkara, A.; Gonzalez, C.; Challacombe, M.; Gill, P. M. W.; Johnson, B.; Chen, W.; Wong, M. W.; Andres, J. L.; Gonzalez, C.; Head-Gordon, M.; Replogle, E. S.; Pople, J. A. *Gaussian 98* (Revision A.7); Gaussian Inc.: Pittsburgh, PA, 1998.
- (27) Schlegel, H. B. *J. Comput. Chem.* **1982**, *3*, 214.
- (28) Becke, A. D. *J. Chem. Phys.* **1993**, *98*, 5648.
- (29) Stephens, P. J.; Devlin, F. J.; Chabalowski, C. F.; Frisch, M. J. *J. Phys. Chem.* **1994**, *98*, 11623.
- (30) Becke, A. D. *Phys. Rev. A* **1988**, *38*, 3098.
- (31) Lee, C. T.; Yang, W. T.; Parr, R. G. *Phys. Rev. B* **1988**, *37*, 785.
- (32) Dunning, T. H. *J. Chem. Phys.* **1989**, *90*, 1007. Woon, D. E.; Dunning, T. H., Jr. *J. Chem. Phys.* **1993**, *98*, 1358.
- (33) Baboul, A. G.; Curtiss, L. A.; Redfern, P. C.; Raghavachari, K. *J. Chem. Phys.* **1999**, *110*, 7650.
- (34) Atkinson, R.; Baulch, D. L.; Cox, R. A.; Hampson, R. F.; Kerr, J. A., Jr.; Troe, J. *J. Phys. Chem. Ref. Data* **1992**, Suppl. No. 4.
- (35) Lias, G.; Bartmess, J. E.; Holmes, J. L.; Levine, R. D.; Mallard, W. G. *J. Phys. Chem. Ref. Data* **1988**, *17*, Suppl. No. 1.
- (36) Kuramshina, G. M.; Weinhold, F.; Pentin, Y. A. *J. Chem. Phys.* **1998**, *109*, 7286.
- (37) Melius, C. F. Sandia National Laboratories, *BAC-MP4 Heats of Formation Data File*, 1997.
- (38) Curtiss, L. A.; Raghavachari, K.; Redfern, P. C.; Pople, J. A. *J. Chem. Phys.* **2000**, *112*, 7374.
- (39) Creasey, J. C.; Lambert, I. R.; Tuckett, R. P.; Codling, K.; Frasinski, L. J.; Hatherly, P. A.; Stankiewicz, M. *J. Chem. Soc., Faraday Trans.* **1991**, *87*, 1287.
- (40) Creasey, J. C.; Lambert, I. R.; Tuckett, R. P.; Codling, K.; Frasinski, L. J.; Hatherly, P. A.; Stankiewicz, M.; Holland, D. M. P. *J. Chem. Phys.* **1990**, *93*, 3295.